

Parametrization of the Driven Betatron Oscillation

R. Miyamoto and S. E. Kopp
Department of Physics
University of Texas at Austin
Austin, Texas 78712 USA

A. Jansson and M. J. Syphers
Fermi National Accelerator Laboratory
Batavia, Illinois 60510 USA
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An AC dipole is a magnet which produces a sinusoidally oscillating dipole field and excites coherent transverse beam motion in a synchrotron. By observing this coherent motion, the optical parameters can be directly measured at the beam position monitor locations. The driven oscillation induced by an AC dipole will generate a phase space ellipse which differs from that of the free oscillation. If not properly accounted for, this difference can lead to a misinterpretation of the actual optical parameters, for instance, of 6% or more in the cases of the Tevatron, RHIC, or LHC. The effect of an AC dipole on the linear optics parameters is identical to that of a thin lens quadrupole. By introducing a new amplitude function to describe this new phase space ellipse, the motion produced by an AC dipole becomes easier to interpret. Beam position data taken under the influence of an AC dipole, with this new interpretation in mind, can lead to more precise measurements of the normal Courant-Snyder parameters. This new parameterization of the driven motion is presented and is used to interpret data taken in the FNAL Tevatron using an AC dipole.

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I. INTRODUCTION

An AC dipole produces a sinusoidally oscillating dipole magnetic field and excites coherent transverse beam motion in a synchrotron for machine diagnosis (Fig 1). Unlike a conventional single turn kicker/pinger magnet, it drives the beam close to the betatron frequency typically for several thousands of revolutions. If the amplitude of its oscillating magnetic field is adiabatically ramped up and down, it can create large coherent oscillations without decoherence and emittance growth [1]. This property makes it a useful diagnosis tool for a proton synchrotron, especially when it is used with an adequate beam position monitor (BPM) system.

AC dipoles have been used in the BNL RHIC [2, 3] and were also tested in the BNL AGS [1] and CERN SPS [4, 5].

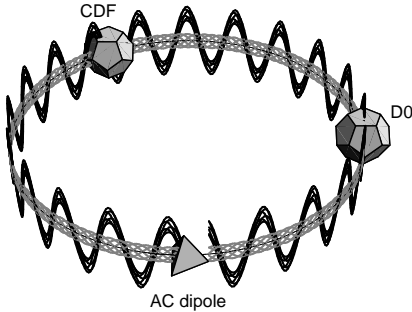


FIG. 1: A diagram of the incoherent free oscillations (gray) and excited coherent oscillations (black) in the Tevatron. Since the oscillations of individual particles within the beam are incoherent, coherent oscillations must be excited to observe betatron motion and measure optical parameters. An AC dipole is a tool to excite sustained coherent transverse oscillations.

Facilitated by its recently upgraded BPM system [6] (now with a resolution of $20 \mu\text{m}$), a vertical AC dipole has been used in the FNAL Tevatron [7, 8, 9]. There is an ongoing project to develop AC dipoles for LHC as well.

When the beam is driven by an AC dipole, the beam motion is governed by two driving terms and the influence of the lesser driving term makes the driven oscillation different from the free oscillation. Although this difference has typically been ignored in previous analyses [3, 10], it could affect the interpretation of the linear optics more than 12% in the Tevatron and 6% in the RHIC and LHC.

This paper proceeds as follows. Section II discusses the two driving terms produced by an AC dipole and presents a new formulation of the driven motion which is suited to treat the two driving terms at the same time. By introducing a new amplitude function β_d for the driven motion, the difference between the free and driven oscillations becomes clear. Section III discusses the difference between the ordinary amplitude function for the free oscillation β and the newly defined amplitude function for the driven oscillation β_d . It is shown that the AC dipole has an analogy with a gradient error and, relative to β , β_d behaves as if there is a gradient error. Section IV presents a few properties of the driven motion which were observed in the Tevatron. The new formulation matches the observed data well.

II. A MODEL OF THE DRIVEN OSCILLATION

A. Two Driving Terms of an Oscillating Dipole Field

The tune of an AC dipole ν_{acd} is defined as the ratio between the frequencies of the AC dipole f_{acd} and the beam

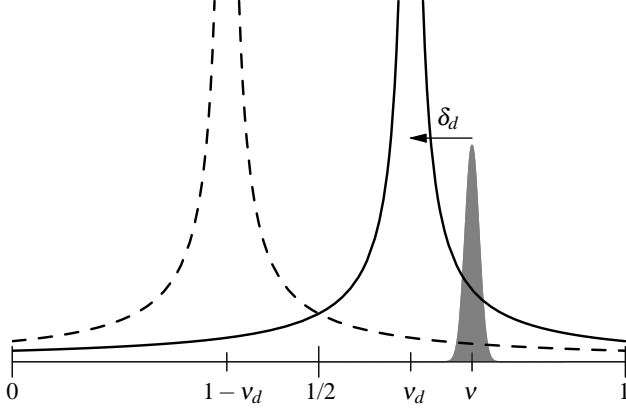


FIG. 2: The amplitude of the driven motion versus the machine tune. A circulating beam is influenced by both (solid and dashed) of the resonant amplitudes. In typical operations of an AC dipole, δ_d is the order of 0.01 so that the primary driving tune ν_d is outside of the tune spread (shaded area).

revolution f_{rev} : $\nu_{\text{acd}} = f_{\text{acd}}/f_{\text{rev}}$. In the following, for any tunes, only their fractional parts are considered. For instance, if $f_{\text{acd}}/f_{\text{rev}}$ is larger than one, ν_{acd} means the fractional part of $f_{\text{acd}}/f_{\text{rev}}$. Since the beam sees an AC dipole only once per revolution, the tune of an AC dipole ν_{acd} and $1 - \nu_{\text{acd}}$ are equivalent (cf. Nyquist sampling theorem). Hence, under the influence of an oscillating dipole field a beam is driven by a pair of driving terms at ν_{acd} and $1 - \nu_{\text{acd}}$. Obviously, the driving term closer to the machine tune ν ($0 < \nu < 1$) has bigger

effects on a beam. In the following, the driving term closer to ν is called the primary and the other is called the secondary. A symbol ν_d is used for the primary driving tune:

$$\nu_d \equiv \begin{cases} \nu_{\text{acd}} & \text{when } |\nu_{\text{acd}} - \nu| < |(1 - \nu_{\text{acd}}) - \nu| \\ 1 - \nu_{\text{acd}} & \text{when } |(1 - \nu_{\text{acd}}) - \nu| < |\nu_{\text{acd}} - \nu| \end{cases} \quad (1)$$

For example, the frequencies of the AC dipole and beam revolution in the Tevatron are $f_{\text{acd}} \simeq 20.5$ kHz and $f_{\text{rev}} \simeq 47.7$ kHz, respectively, and hence the tune of the AC dipole is $\nu_{\text{acd}} = 20.5/47.7 \simeq 0.43$. Since the machine tune of the Tevatron is $\nu \simeq 0.58$, $1 - \nu_{\text{acd}} \simeq 0.57$ is the primary driving tune and $\nu_{\text{acd}} \simeq 0.43$ is secondary in this case (Table I).

The distance from the primary driving term to the machine tune $\delta_d \equiv \nu_d - \nu$ is an important parameter of the driven betatron oscillation. As seen later, the secondary driving term generates a difference between the free and driven oscillations and affects linear optics measurements. Ideally, if the beam is driven very close to the machine tune ν ($\delta_d \rightarrow 0$), the influence of the primary driving term becomes dominant and the secondary driving term can be ignored. In reality, however, the finite tune spread of the beam causes beam losses if $|\delta_d|$ is too small and there is always a lower limit for $|\delta_d|$ (Fig 2). AC dipoles are currently used in the Tevatron and RHIC and planned for the LHC. In these synchrotrons, the lower limit of $|\delta_d|$ is about 0.01 to prevent beam losses.

When the amplitude of the field is constant, the position of the driven beam x_d is given by [10, 11]

$$x_d(nC + \Delta s) \simeq \frac{\theta_{\text{acd}} \sqrt{\beta_{\text{acd}}}}{4 \sin[\pi(\nu_{\text{acd}} - \nu)]} \sqrt{\beta(\Delta s)} \cos[2\pi\nu_{\text{acd}}n + \psi(\Delta s) + \pi(\nu_{\text{acd}} - \nu) + \chi_{\text{acd}}] \\ + \frac{\theta_{\text{acd}} \sqrt{\beta_{\text{acd}}}}{4 \sin[\pi((1 - \nu_{\text{acd}}) - \nu)]} \sqrt{\beta(\Delta s)} \cos[2\pi(1 - \nu_{\text{acd}})n + \psi(\Delta s) + \pi((1 - \nu_{\text{acd}}) - \nu) - \chi_{\text{acd}}], \quad (2)$$

where C is the circumference of a ring, Δs ($0 \leq \Delta s < C$) is the longitudinal position measured from the location of the AC dipole, θ_{acd} is the maximum kick angle of the AC dipole, β_{acd} is the amplitude function at the location of the AC dipole, ψ is the phase advance of the free oscillation measured from the location of the AC dipole, and χ_{acd} is the initial phase of the AC dipole. The two terms in Eq 2 are completely symmetric and represent the influences of the two driving terms [14]. To quantify the effect of the secondary driving term, it is useful to define a parameter which describes the ratio between the primary (larger) and secondary (smaller) modes in Eq 2:

$$\lambda_d(\delta_d) \equiv \frac{\sin[\pi(\nu_d - \nu)]}{\sin[\pi((1 - \nu_d) - \nu)]} = \frac{\sin(\pi\delta_d)}{\sin(2\pi\nu + \pi\delta_d)}. \quad (3)$$

When $|\delta_d| = 0.01$, $|\lambda_d| \simeq 0.06$ for the Tevatron with $\nu \simeq 0.58$ and $|\lambda_d| \simeq 0.03$ for the RHIC and LHC with $\nu \simeq 0.3$ and 0.7 (Table I). This is the effect of the secondary driving term on

the amplitude of the driven oscillation. When the machine tune is closer to the half-integer, the two driving terms are closer to each other and the influence of the secondary driving term gets larger. This is why $|\lambda_d|$ of the Tevatron is larger than that of the RHIC and LHC.

TABLE I: Parameters related to the driven oscillation in the Tevatron, RHIC, and LHC when $|\delta_d| = 0.01$. The secondary driving term affects the amplitude of the driven motion by $|\lambda_d|$ and, as seen later, produces an effect like β -beat with the amplitude of $2|\lambda_d|$.

Parameter	Tevatron	RHIC	LHC
Machine Tune ν	.58	.7	.3
AC Dipole Tune ν_{acd}	.42	.7	
$ \lambda_d $	6%	3%	3%
Amplitude of the β -beat	12-13%	6-7%	6-7%

B. A New Parametrization of the Driven Betatron Oscillation

Eq 2 can be written in the following compact form which includes the influences of both driving terms:

$$x_d(s; \delta_d) = A_d(\delta_d) \sqrt{\beta_d(s; \delta_d)} \cos(\psi_d(s; \delta_d) \pm \chi_{\text{acd}}). \quad (4)$$

Here, A_d is a quantity with dimensions of (length)^{1/2}:

$$A_d(\delta_d) = \frac{\theta_{\text{acd}}}{4 \sin(\pi \delta_d)} \sqrt{(1 - \lambda_d(\delta_d)^2) \beta_{\text{acd}}}, \quad (5)$$

β_d is a newly defined amplitude function of the driven oscillation which satisfies

$$\frac{\beta_d(s; \delta_d)}{\beta(s)} = \frac{1 + \lambda_d(\delta_d)^2 - 2\lambda_d(\delta_d) \cos(2\psi(s) - 2\pi\nu)}{1 - \lambda_d(\delta_d)^2}, \quad (6)$$

ψ_d is a newly defined phase advance of the driven oscillation measured from the location of the AC dipole:

$$\psi_d(s; \delta_d) = \int_0^s \frac{d\bar{s}}{\beta_d(\bar{s}; \delta_d)}, \quad (7)$$

and the sign in front of χ_{acd} is positive when $\nu_d = \nu_{\text{acd}}$ and negative when $\nu_d = 1 - \nu_{\text{acd}}$. Hence, the driven oscillation can be parametrized in the same form as the free oscillation even when the influences of the both driving terms are included. Since A_d is a constant of motion, the difference between the free and driven oscillations comes from the amplitude function β_d and phase advance ψ_d . As discussed previously, in the limit, $\nu_d \rightarrow \nu$, the primary driving term becomes dominant and the secondary driving term can be ignored. In this limit $\lambda_d \rightarrow 0$ and β_d and ψ_d converge to β and ψ .

If the lesser mode in Eq 2 is ignored, the oscillation phase has an apparent jump of $2\pi\delta_d$ at the location of the AC dipole. However, if the influences of both driving terms are properly included as Eq 4, the phase advance is smooth at the location of the AC dipole. A relation between the phase advances of free and driven oscillations, ψ and ψ_d , is given by

$$\begin{aligned} \tan(\psi_d - \pi\nu_d) &= \frac{1 + \lambda_d}{1 - \lambda_d} \tan(\psi - \pi\nu) \\ &= \frac{\tan(\pi\nu_d)}{\tan(\pi\nu)} \tan(\psi - \pi\nu). \end{aligned} \quad (8)$$

For the free oscillation, the phase advance in a single revolution is $\psi(s+C) - \psi(s) = 2\pi\nu \pmod{2\pi}$. In the equation above, $\psi_d = 2\pi\nu_d$ when $\psi = 2\pi\nu$. Hence, the phase advance in a single revolution is $2\pi\nu_d$ for the driven motion.

III. DIFFERENCE BETWEEN THE AMPLITUDE FUNCTIONS β AND β_d

As seen in the previous section, the difference between the free and driven oscillations lies in the difference of their amplitude functions, β and β_d . It is crucial to understand this difference between β and β_d in detail when an AC dipole is used to diagnose a synchrotron.

In free betatron oscillations tune and amplitude function are coupled, and a change in tune involves a change in amplitude function and vice versa. This is true for the driven betatron oscillation, too. As seen in the previous section, for the driven oscillation, both the amplitude function and tune, β_d and ν_d , are different from those for the free oscillation, β and ν . As a matter of fact, the relation between these changes of the tune and amplitude function is the same as that for a gradient error. Hence, reviewing the effect of a gradient error is helpful to understand the driven oscillation.

A. Review of a Gradient Error

If a synchrotron has a gradient error, its machine tune ν and amplitude function β change [12]. Suppose a synchrotron has a gradient error with the strength $q_{\text{err}} = B'\ell/(B\rho)$ at the longitudinal position $s = 0$. Then, the equation of motion is given by

$$x'' + k(s)x = -q_{\text{err}} \left[\sum_{n=-\infty}^{\infty} \delta(s - Cn) \right] x, \quad (9)$$

where the prime denotes the derivative with the longitudinal coordinate s , k is the spring constant, and δ is the Dirac's delta function.

By comparing the single turn transfer matrices with and without the gradient error, the new tune ν_q and amplitude function β_q satisfy the following two equations [12]:

$$q_{\text{err}} = 2 \frac{\cos(2\pi\nu) - \cos(2\pi\nu_q)}{\beta_{\text{err}} \sin(2\pi\nu)} \quad (10)$$

$$\frac{\beta_q}{\beta} = \frac{\sin(2\pi\nu)}{\sin(2\pi\nu_q)} - q_{\text{err}} \beta_{\text{err}} \frac{\sin \psi \sin(2\pi\nu - \psi)}{\sin(2\pi\nu_q)}, \quad (11)$$

where β_{err} is the amplitude function at the gradient error and ψ is the phase advance measured from the gradient error. By substituting the first equation into the second, the ratio between the new and original amplitude functions β_q/β is given by

$$\frac{\beta_q}{\beta} = \frac{1 + \lambda_q^2 - 2\lambda_q \cos(2\psi - 2\pi\nu)}{1 - \lambda_q^2}. \quad (12)$$

Here, λ_q is a parameter similar to λ_d in Eq 3:

$$\lambda_q \equiv \frac{\sin(\pi\delta_q)}{\sin(2\pi\nu + \pi\delta_q)}, \quad (13)$$

where δ_q is the tune shift by a gradient error $\delta_q \equiv \nu_q - \nu$. When the gradient error q_{err} is small, the new and original amplitude functions satisfy

$$\frac{\beta_q - \beta}{\beta} \simeq -2\lambda_q \cos(2\psi - 2\pi\nu). \quad (14)$$

This quantity behaves like a standing wave in a synchrotron and is called the β -beat (or sometimes β -wave). The amplitude of the β -beat is $2|\lambda_q|$.

B. Analogy to a Gradient Error

As seen in Eqs 6 and 12, the relation between β_d and δ_d for an oscillating dipole field is the same as the relation between β_q and δ_q for a gradient error. The following argument gives insight why an oscillating dipole field changes the observed phase space motion as like a gradient error.

When the oscillation amplitude of the AC dipole field is constant, the Hill's equation of motion is given by

$$x'' + k(s)x = - \sum_n \theta_{\text{acd}} \cos(2\pi\nu_d n \pm \chi_{\text{acd}}) \delta(s - Cn). \quad (15)$$

The right-hand-side describes the kicks by the AC dipole located at $s = 0$. The summation runs over the time period when the oscillation amplitude of the AC dipole field is constant and the sign in front of the initial phase χ_{acd} is the same convention as Eq 4. Eq 4 is the particular solution of this inhomogeneous Hill's equation when the oscillation amplitude of the AC dipole field is adiabatically ramped to a constant amplitude. Since the phase of the driven oscillation ψ_d increases by $2\pi\nu_d \pmod{2\pi}$ in one revolution, the position of the driven oscillation at the location of the AC dipole $s = Cn$ is given by

$$x_d(Cn; \delta_d) = A_d(\delta_d) \sqrt{\beta_d(0; \delta_d)} \cos(2\pi\nu_d n \pm \chi_{\text{acd}}). \quad (16)$$

Notice the phases of the driven oscillation x_d and the kicks by the AC dipole in Eq 15 are both $2\pi\nu_d n \pm \chi_{\text{acd}}$ at the location of the AC dipole. Hence, when the beam passes the AC dipole, its magnetic field is proportional to the position of the driven oscillation x_d like a quadrupole magnet. This is the physical reason why an oscillating dipole field changes the amplitude function like a gradient error. The phases of the driven oscillation and the AC dipole are synchronized like this only when the oscillation amplitude of the AC dipole field is constant after the adiabatic ramp up. Since x_d is the solution of Eq 15, it formally satisfies the following equation

$$x_d'' + k(s)x_d = -q_{\text{acd}} \left[\sum_n \delta(s - Cn) \right] x_d. \quad (17)$$

Here, Eq 16 is used to change the right-hand-side and q_{acd} is a constant given by

$$q_{\text{acd}} = \frac{\theta_{\text{acd}}}{A_d \sqrt{\beta_d(0; \delta_d)}} = 2 \frac{\cos(2\pi\nu) - \cos(2\pi\nu_d)}{\beta_{\text{acd}} \sin(2\pi\nu)}. \quad (18)$$

Eq 17 is exactly the same as the Hill's equation with a gradient error, Eq 9. By comparing Eqs 9, 10, 18, and 17, it is trivial that the relation between β_d and δ_d is the same as the relation between β_q and δ_q .

C. Ring-wide Behavior of β_d

As discussed in the previous two sections, for the driven motion, the observed amplitude function β_d differs from the actual β as if there is a gradient error. Hence, β_d is beating relative to β and the beating amplitude is about $2|\lambda_d|$ from

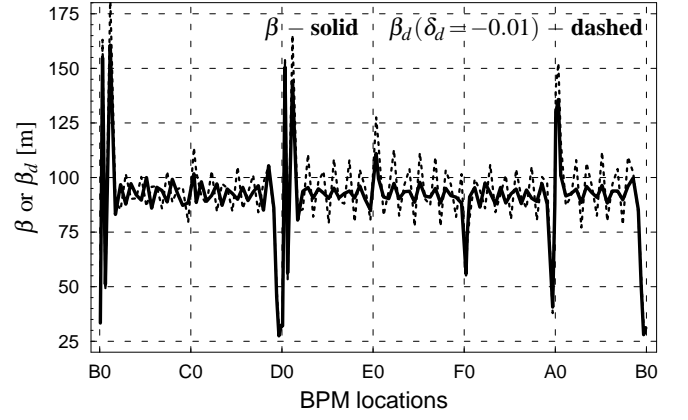


FIG. 3: The amplitude functions of the free and driven oscillations, β (solid) and β_d when $\delta_d = -0.01$ (dashed). Both of them are calculated from turn-by-turn data of the driven oscillation. As expected, β_d shows the 10-15% beating relative to β . If the difference of β and β_d is simply ignored, the β measurement has this much error and the real β -beat cannot be distinguished from the beating of β_d .

Eq 14. Remember the effect of the secondary driving term on the beam motion is the order of λ_d . Since the amplitude function is proportional to the square of the position, its effect on the amplitude function is of the order of $2|\lambda_d|$. Since the minimum difference between the primary driving tune and machine tune $|\delta_d|$ is about 0.01 for the Tevatron, RHIC, and LHC, the beating amplitude of β_d relative to β is 12-13% for the Tevatron and 6-7% for the RHIC and LHC (Table I).

When turn-by-turn beam positions at all BPMs are given for the free oscillation, the relative β -function can be determined by simply comparing the square of the oscillation amplitude at each BPM. If the same analysis is applied to the turn-by-turn data of the driven oscillation, what is calculated is β_d instead of β . If the difference between β_d and β is simply ignored and β is determined in this way, the error may be as large as $2|\lambda_d|$. Furthermore, since the beating of β_d cannot be distinguished from the real β -beat caused by gradient errors, the real β -beat cannot be measured in this way without depending on a machine model.

To calculate the true β -function from turn-by-turn data of the driven oscillation without depending on a machine model, multiple sets of data are necessary [8]. Fig 3 shows amplitude functions of the free and driven oscillations, β and $\beta_d(\delta_d = -0.01)$. They are both measured from data of the driven oscillation. Multiple data sets are used to calculate β as described in [8] and β_d is calculated by comparing the square of the amplitude at each BPM. As expected, β_d is showing the beating of 10-15% relative to β .

D. Relation between β_d and δ_d

The previous section discussed the global behavior of β_d compared to β . This section considers how β_d changes depending on δ_d at one location of a synchrotron. From Eq 6, the relation between β_d and β becomes nonlinear when λ_d is

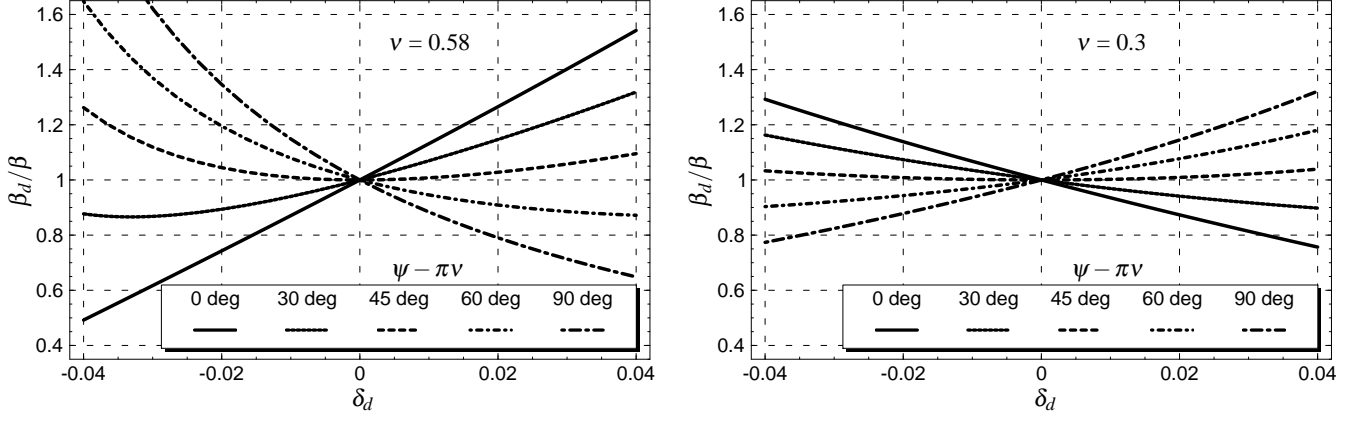


FIG. 4: The relation between the amplitude functions of the free and driven betatron oscillations β and β_d . The ratio β_d/β is numerically calculated based on Eq 6 by changing the difference between the primary driving tune and machine tune δ_d and the phase advance ψ . The left plot is when the machine tune 0.58 like the Tevatron and the right is when 0.3 like the RHIC and LHC. Since the secondary driving term gets closer and λ_d gets larger when the machine tune is closer to the half integer, β_d/β is larger and the nonlinearity is stronger in the left plot. The nonlinearity gets larger when $\psi - \pi\nu$ gets closer to 45 deg and $\cos(2\psi - 2\pi\nu)$ gets closer to zero.

large or the phase term $\cos(2\psi - 2\pi\nu)$ is close to zero. Since the difference between β and β_d has a considerable impact on the linear optics measurement, it is important to understand the properties of Eq 6 over wide ranges of parameters. Fig 4 shows the numerical calculations of β_d/β based on Eq 6. The two plots are for two different machine tunes: $\nu = 0.58$ like the Tevatron and $\nu = 0.3$ like the RHIC and LHC. Since λ_d is almost twice as large for the same δ_d when $\nu = 0.58$ compared to $\nu = 0.3$, the nonlinearity grows much faster with δ_d in the Tevatron. It is also seen in the left plot that the nonlinearity becomes larger when $|\cos(2\psi - 2\pi\nu)|$ gets closer to zero. Such a nonlinear relation between β_d and δ_d can be actually seen for the driven oscillation excited in the Tevatron. An example is shown in the next section.

IV. EVIDENCE OF THE SECONDARY DRIVING TERM

A. Rotation of the Phase Space Ellipse

The previous section discussed the amplitude function of the driven motion β_d . Parameters corresponding to the other Courant-Snyder parameters α and γ can be also defined as for the free oscillation:

$$\alpha_d(s; \delta_d) \equiv -\frac{1}{2} \frac{d\beta_d(s; \delta_d)}{ds} \quad (19)$$

$$\gamma_d(s; \delta_d) \equiv \frac{1 + \alpha_d(s; \delta_d)^2}{\beta_d(s; \delta_d)}. \quad (20)$$

The explicit forms of these parameters are given by

$$\alpha_d = \frac{1 + \lambda_d^2 - 2\lambda_d \cos(2\psi - 2\pi\nu)}{1 - \lambda_d^2} \alpha - \frac{2\lambda_d \sin(2\psi - 2\pi\nu)}{1 - \lambda_d^2} \quad (21)$$

and

$$\gamma_d = \frac{1 + \lambda_d^2 + 2\lambda_d \cos(2\psi - 2\pi\nu + 2 \arctan \alpha)}{1 - \lambda_d^2} \gamma. \quad (22)$$

When β_d , α_d , γ_d , and A_d are defined this way, they satisfy the Courant-Snyder invariance:

$$A_d^2 = \gamma_d x_d'^2 + 2\alpha_d x_d x_d' + \beta_d x_d'^2. \quad (23)$$

Hence, the turn-by-turn position and angle of the driven oscillation also form an ellipse on the phase space, like the free oscillation. Since not only A_d but also the Courant-Snyder-like parameters β_d , α_d , and γ_d depend on the difference between the primary driving tune and the machine tune δ_d , both

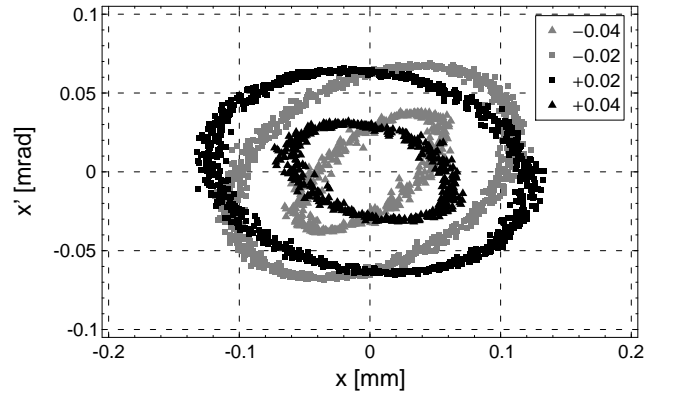


FIG. 5: Measured phase space ellipses of the driven oscillations when $\delta_d = \pm 0.02$ and ± 0.04 . Here, δ_d is the difference between the primary driving tune and the machine tune. The figure shows the phase space at one of the low- β locations (B0) in the Tevatron where the derivative of the amplitude function α is zero by design. Since the Courant-Snyder-like parameters of the driven oscillation β_d , α_d , and γ_d depend on δ_d , not only the area but also the shape of the ellipse changes with δ_d .

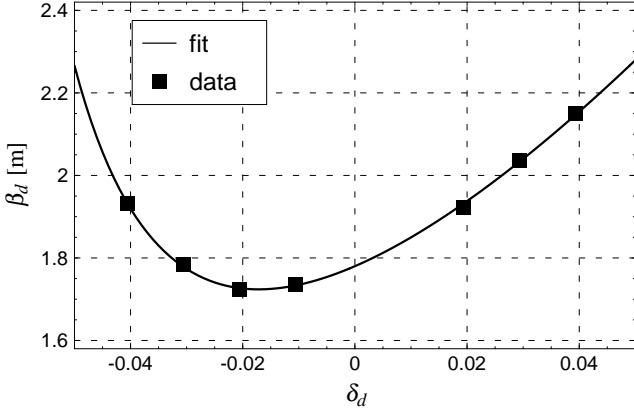


FIG. 6: The relation between the amplitude function of the driven oscillation β_d and the difference between the primary driving tune and the machine tune δ_d . The location is the same low- β point of the Tevatron (B0) in Fig 5. The amplitude function β_d at each data point is determined from the shape of an ellipse in Fig 5. The curve is the fit of Eq 6 to the data points. Despite the strong nonlinearity, Eq 6 is fitting well. In the figure, $\beta_d(\delta_d = 0)$ corresponds to the true value of β at the location.

the area and shape of the phase space ellipse changes with δ_d for the driven oscillation. In two collision straight sections of the Tevatron, B0 and D0, there are pairs of BPMs with no magnetic element in-between. The beam travels along straight lines between these pairs and, hence, both position and angle can be directly measured at these locations. Fig 5 shows the measured phase ellipses of the driven oscillations by using a pair of such BPMs. The frequency of the AC dipole was changed to adjust δ_d to ± 0.04 and ± 0.02 , while the kick angle of the AC dipole θ_{acd} was kept the same. As expected, the shape of the phase space ellipse changes with δ_d . Since δ_d dependence of β_d , α_d , and γ_d comes from the secondary driving term, the rotation of the phase space ellipse is its qualitative evidence.

By fitting Eq (23) to an ellipse in Fig 5, its area πA_d^2 and the parameters β_d , α_d , and γ_d can be determined. Fig 6 shows β_d determined from the fits to ellipses in Fig 5 (and three more). The curve in the figure is the fit of Eq 6 to the data with parameters β and ψ . The model of Eq 6 is fitting well to the data even though the nonlinearity is strong in the relation between β_d and δ_d at the location. The β -function at the location can be calculated as one of the fit parameters. In the figure it is the value of β_d when $\delta_d = 0$.

B. Asymmetric Amplitude Response

When the secondary driving term is negligible, by ignoring the smaller term of Eq 2 or taking the limit of $\lambda_d \rightarrow 0$ in Eqs 4, 5, and 6, the amplitude of the driven oscillation can be approximated by

$$a_d^{(0)}(s; \delta_d) \equiv \frac{\theta_{acd} \sqrt{\beta_{acd} \beta}}{4 |\sin(\pi \delta_d)|}. \quad (24)$$

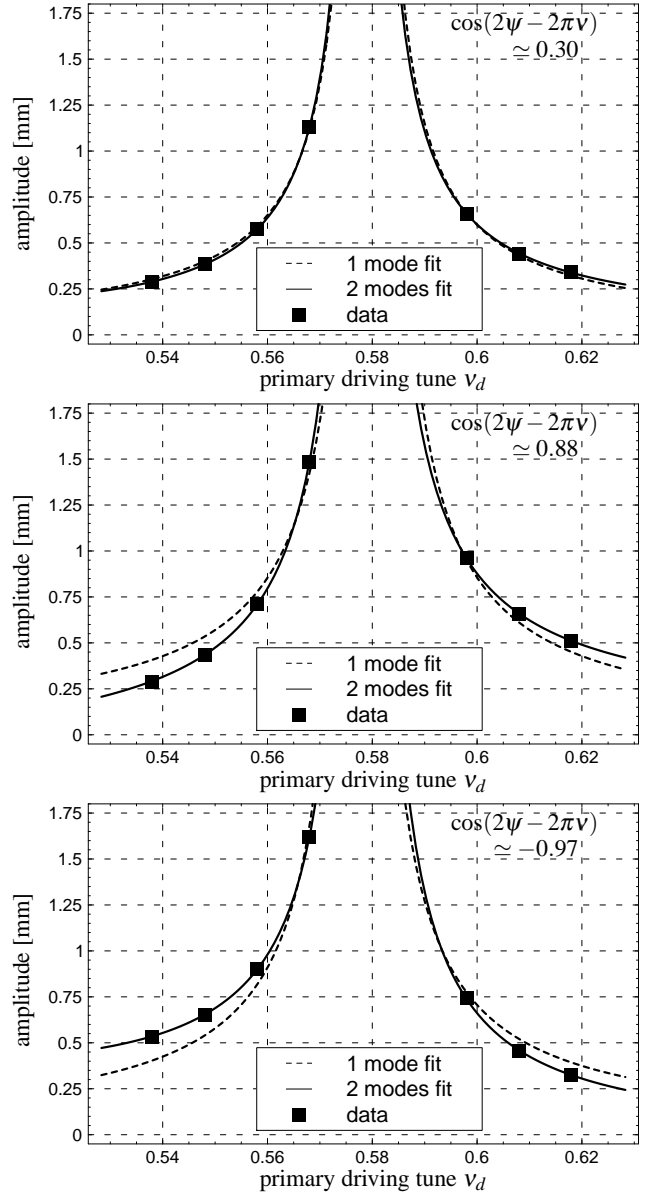


FIG. 7: Measured relation between the amplitude of the driven oscillation and the primary driving tune ν_d at three BPM locations in the Tevatron. In the measurements, only ν_d was changed. The solid and dashed lines are fits with and without the effect of the secondary driving term. The asymmetry around the tune $\nu \approx 0.5785$ increases with $|\cos(2\psi - 2\pi\nu)|$. In the second plot where $\cos(2\psi - 2\pi\nu) > 0$, the amplitude is larger in the region $\nu_d > \nu$. As expected from Eq 26, the relation flips in the third plot where $\cos(2\psi - 2\pi\nu) < 0$.

In this case, the amplitude of the driven oscillation depends on the primary driving tune ν_d only through $|\sin(\pi \delta_d)|$ (remember $\delta_d = \nu_d - \nu$) and is symmetric around the machine tune ν . From Eqs 4, 5, and 6, the amplitude including the effect of the secondary driving term $a_d(s; \delta_d)$ is given by

$$a_d(s; \delta_d) = a_d^{(0)} \sqrt{1 + \lambda_d^2 - 2\lambda_d \cos(2\psi - 2\pi\nu)}. \quad (25)$$

Now, the amplitude depends on v_d through the factor $[1 + \lambda_d^2 - 2\lambda_d \cos(2\psi(s) - 2\pi v)]^{1/2}$ as well. To the first order in δ_d , the amplitude is approximated by

$$a_d \simeq a_d^{(0)} \left[1 - \frac{\pi \cos(2\psi(s) - 2\pi v)}{\sin(2\pi v)} \delta_d \right]. \quad (26)$$

Hence, the secondary driving term makes the v_d dependence of the amplitude asymmetric around the machine tune v . The magnitude of this asymmetry at each location is determined by the factor $\cos(2\psi - 2\pi v)$.

Fig 7 shows the relation between the amplitude of the driven oscillation and v_d at three BPM locations in the Tevatron. In the measurements, only the frequency of the AC dipole was changed while its kick angle θ_{acd} was kept the same. The dashed and solid lines represent the fits of Eq 24 and Eq 25 to the data. The fit parameters are $\theta_{acd}[\beta_{acd}\beta]^{1/2}$ and v for Eq 24 and $\theta_{acd}[\beta_{acd}\beta]^{1/2}$, v , and ψ for Eq 25 [15]. At two locations where $|\cos(2\psi - 2\pi v)|$ is close to one, the asymmetry around the machine tune ($v \simeq 0.5785$) is large and the result of the fits without the secondary driving term (Eq 24) is not well matched.

Although the existence of a secondary driven term effect is clear in Fig 7, there is better evidence that Eq 25 fits the data better than 24. From the fits in Fig 7, the machine tune v can be determined at each BPM location. Fig 8 shows machine tunes determined at all BPM locations from the fits of the amplitude versus v_d . The dashed and solid lines represent machine tunes from the fits of Eqs 24 and 25. Since the machine tune v is a global parameter of a synchrotron, the variation of the measured machine tune over BPMs shows the inaccuracy of the measurement. From the figure, it is clear the model including the secondary driving fits to the data much better. This also shows the importance of the secondary driving term in the driven oscillation.

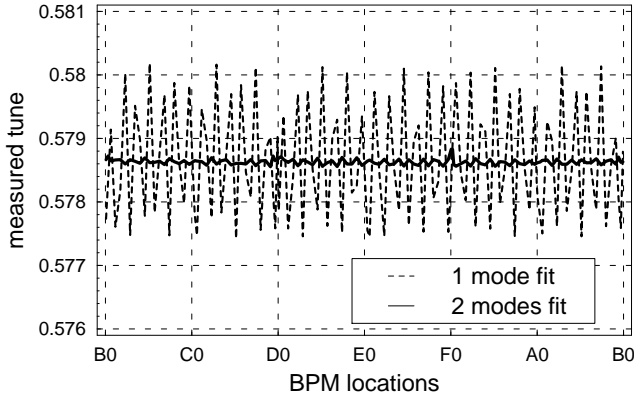


FIG. 8: Measured machine tune at all BPM locations from the fits of the amplitude versus the primary driving tune in Fig 7. The solid line includes the influence of the secondary driving term and the dashed line does not. Since the tune is a global parameter of a synchrotron, the variation of the measured machine tune over BPMs shows the inaccuracy of the measurement. From the figure, it is clear the model including the secondary driving fits to the data much better. This also shows the importance of the secondary driving term in the driven oscillation. The beating of the dashed line is caused by ignoring the beating of β_d relative to β .

V. CONCLUSION

Under the influence of a sinusoidally oscillating magnetic field of an AC dipole, the beam is driven by two driving terms. As a result the phase space trajectory of the driven motion is different from that of the free betatron motion. If the difference is simply ignored, interpretation of the linear optics based on the data of the driven motion can have 12-13% error for the Tevatron and 6-7% error for the RHIC and LHC. This difference on the phase space is identical to the influence of a gradient error at the same location as the AC dipole. Hence, as a gradient error changes the amplitude function around the ring, the expression of the driven motion can be simplified by introducing a new amplitude function for the driven motion.

This paper presented a few examples of the difference between the free and driven motions as observed in the Tevatron. It also showed that the new parametrization of the driven motion clarifies the data interpretation and multiple data sets are necessary to better resolve the true optical parameters of the free oscillation.

With this knowledge, very precise and direct measurements of the amplitude function in a hadron synchrotron can be obtained quickly without degradation of the beam quality, using a small number of data sets obtained at different frequencies of the AC dipole. This technique will be especially useful in the LHC, for example, to adjust the beam envelope at critical locations such as at beam collimation devices.

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 - [14] Eq 2 is assuming the amplitude of the AC dipole field is adiabatically ramped up to a constant value. The exact expression of x_d includes transient modes which are inversely proportional to the ramp up time and oscillate with the machine tune ν . If the ramp up is slow enough, all of these modes are very small and decohere before the end of the ramp up. Hence these ignored modes do not affect the motion of the beam centroid but they may affect the beam size [11].
 - [15] The ring-wide β -function can be determined from the fit up to a constant $\theta_{\text{acd}} \sqrt{\beta_{\text{acd}}}$. The constant is determined from the analysis using a pair of BPMs in the collision straight sections. See [8] for details. The phase advance ψ can be determined for the fit, too.